

Steady-state simulation of homogeneous vapor-liquid nucleation "by the intervention of intelligent beings"

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MD simulation of a single droplet

- Vapor and liquid are equilibrated separately
- A small (j < 10000) droplet is inserted into the vapor
- If the droplet cannot evaporate completely, an equilibrium is established within a few nanoseconds



t-s-LJ fluid (r_c = 2.5 σ)





Single droplet in non-equilibrium

t-s-LJ fluid 200 16000 Vapor properties: 2000 150 N = 130,000100 $\rho = 0.0268 / \sigma^3$ 1000 $T = 0.80 \ \epsilon/k$ 50 Time interval \approx 1 ns 0 Classical theory: $j^* = 850$ 200 400 600 800

The critical nucleus size cannot be clearly detected by non-equilibrium simulation





Equilibrium vapor pressure



Equilibrium condition for a droplet containing *j* molecules:

$$p = p(T, j)$$

 ΔG at constant *p* and *T*: 1 unstable equilibrium

 ΔF at constant V and T:

1 unstable equilibrium 1 stable equilibrium





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Critical nucleus from equilibrium MD simulation



Kelvin equation:

$$j^* = \left(\frac{2\gamma A_1}{3\Delta \mu}\right)^3$$

- present simulation data
- classical theory
- Good prediction of the critical size for low temperatures
- Above 0.8 T_c, significant deviations are present





Nucleation from MD in the *NVT* ensemble

Yasuoka and Matsumoto (1998):

Number of emerging nuclei with >*j* molecules per volume and time

Typical approach:

determine the nucleation rate for various *j* values

250,000 methane molecules at 130 K and 1.606 mol/l (full LJ potential)







Nucleation rate following Yasuoka & Matsumoto







Grand canonical MD

Algorithm: MC insertion/deletion steps alternating with MD steps

- fixed values of μ , V und T
- test insertion of a molecule at a random position

$$\boldsymbol{P}_{\text{ins}} = \min\left[1, \exp\left(\frac{\mu - \Delta \boldsymbol{U}_{\text{pot}}}{kT}\right) \frac{V}{\Lambda^{3} (N+1)}\right]$$

• test deletion of a random molecule

$$\boldsymbol{P}_{del} = \min\left[1, \exp\left(\frac{-\mu - \Delta \boldsymbol{U}_{pot}}{kT}\right) \frac{\boldsymbol{V}}{\Lambda^3 \boldsymbol{N}}\right]$$

• equal number of insertions and deletions $(10^{-6}N \text{ to } 10^{-3}N / \text{ step})$

Thermodynamic conditions of supersaturated state are maintained





McDonald's demon







Intervention rate as a function of the threshold size

In J j*_{CNT} t-s-LJ fluid $J_i = J/q_i$ $T = 0.7 \epsilon k$ S = 2.496-15 Probability for a nucleus of growing from size *i* to infinite size: CNT shifted $q_{i} = \frac{\int_{1}^{\prime} \exp(2\beta\Omega) \, dj}{\int_{1}^{\infty} \exp(2\beta\Omega) \, dj},$ -20 to $\ln J = -22.0$ Δ Δ -25 Δ CNT $q_{i^*} \approx \frac{1}{2}$ and in particular: i 20 60 0 40

CNT predicts an acceptable value for j* and underestimates J significantly





IhEt

Maximal overheating of growing nuclei



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GCMD simulation of nucleation: Results







Conclusion

- MD simulation of steady states allows sampling over an arbitrary time interval, eventually leading to the desired level of accuracy
- Single droplets can be stable in the canonical ensemble
- A supersaturated vapor near the spinodal line can be stabilized by grand canonical simulation with McDonald's demon
- Classical nucleation theory leads to acceptable results for the t-s-LJ fluid

