

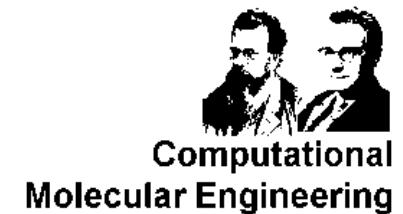


# Non-equilibrium molecular dynamics simulation of real fluids in nanoporous materials

Martin Horsch, Hendrik Frentrup, Carlos Avendaño Jiménez,  
Alicia Marín Torres, Alaaeldin Salih, Jadran Vrabec, Erich Müller, Hans Hasse

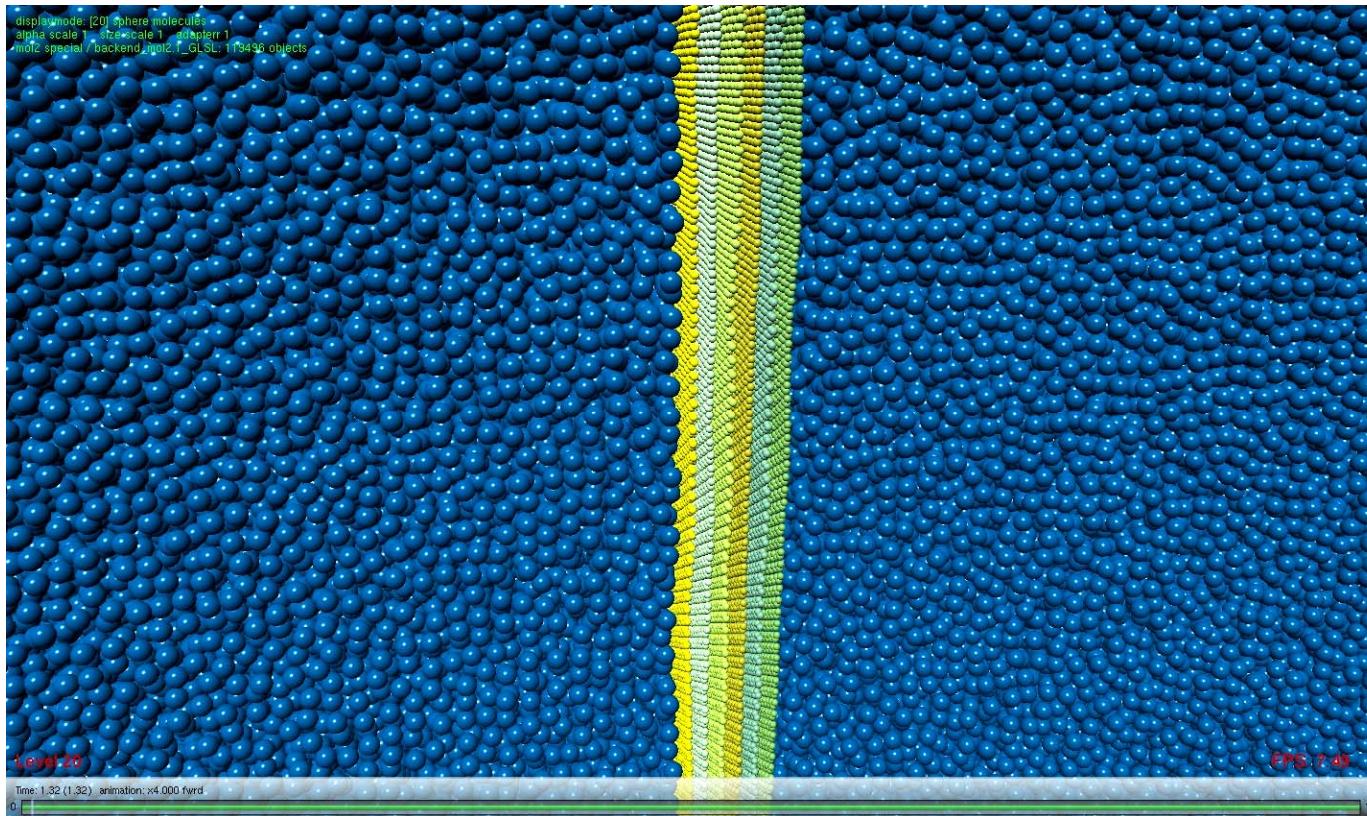
<sup>1</sup>TU Kaiserslautern, <sup>2</sup>Imperial College London, <sup>3</sup>Cornell U., <sup>4</sup>U. Paderborn

Workshop „Industrial Use of Molecular Thermodynamics“,  
Lyon, 19<sup>th</sup> March 12



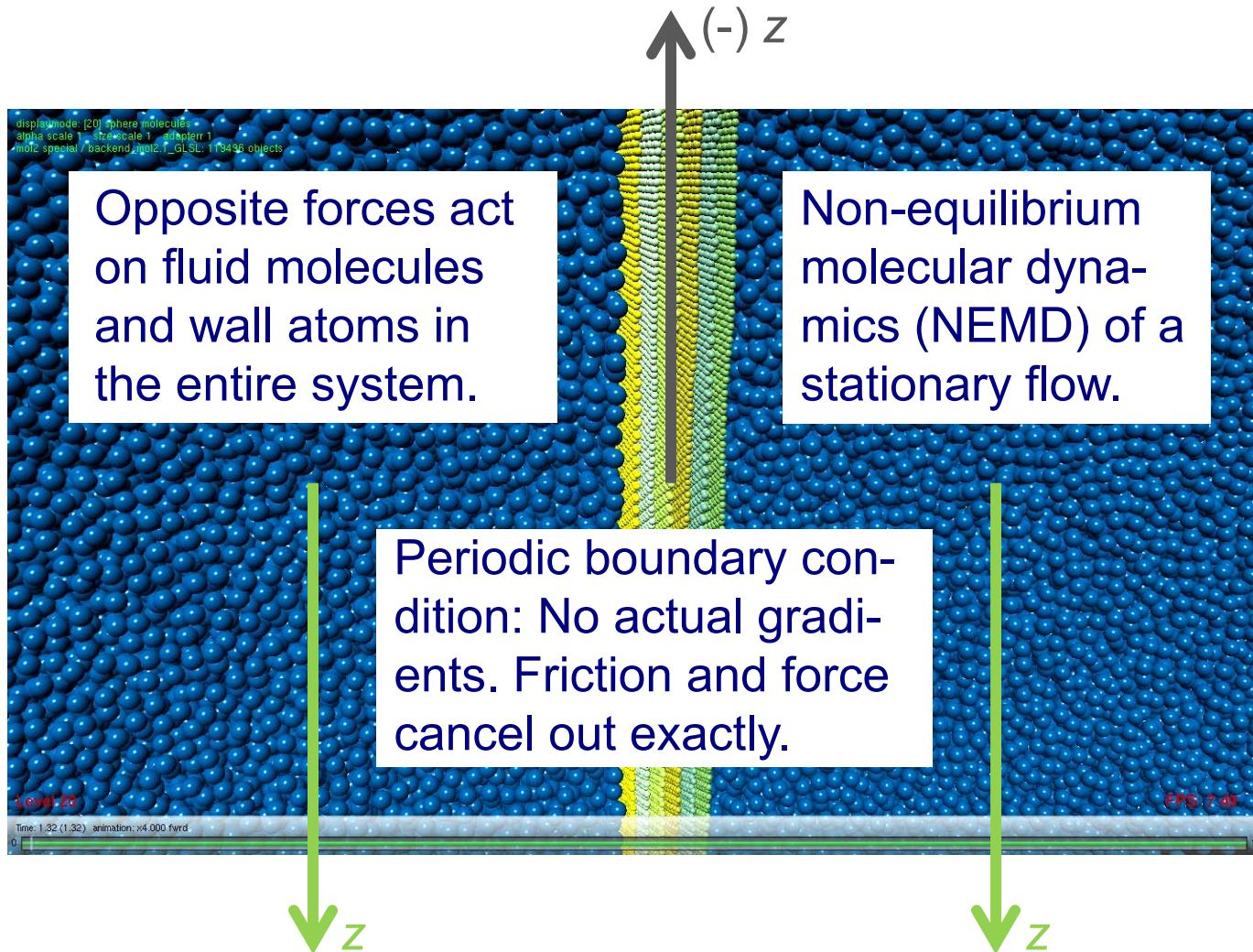


# Exact compensation of the pressure drop



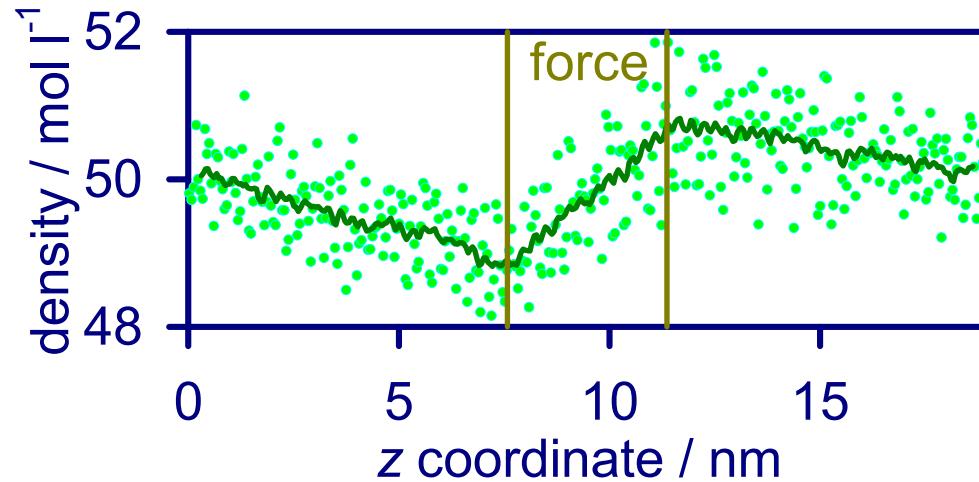


# Exact compensation of the pressure drop



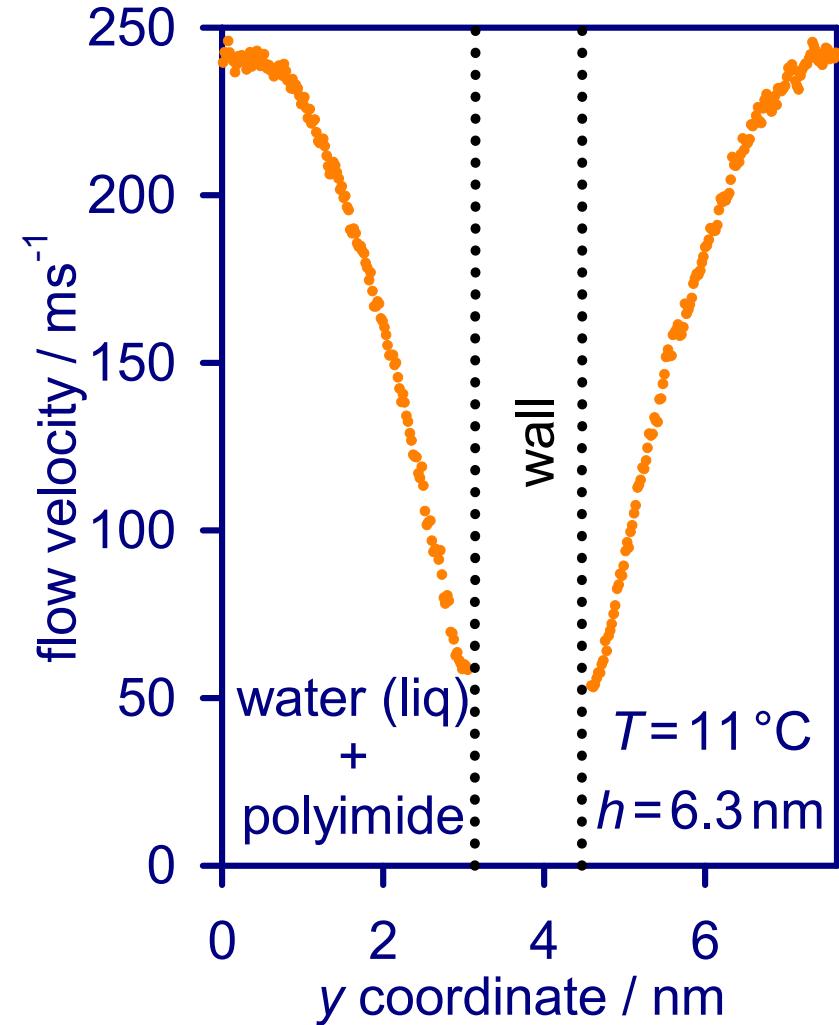


# Overcompensation of the pressure drop



The accelerating force is only applied to the fluid molecules within a specified control volume.

It overcompensates the pressure drop, so that (equivalent) density, pressure, and chemical potential gradients are actually present.





# Viscous and diffusive mass transfer

The transport diffusion coefficient  $D_t$  consists of contributions from two different mechanisms:

$$J_i = -D_t \nabla \rho_i = -L_f \nabla (\mu_i / T) - D'_t \nabla p$$

**Diffusive transport:** Mobility  $D_0 = D_s + D_\xi$ , expressed above in terms of the Onsager type coefficient  $L_f$ , caused by the random thermal movement of individual molecules.

**Viscous transport:** Ordered collective motion of the molecules due to a pressure gradient, which can be understood in terms of a fluid continuum.

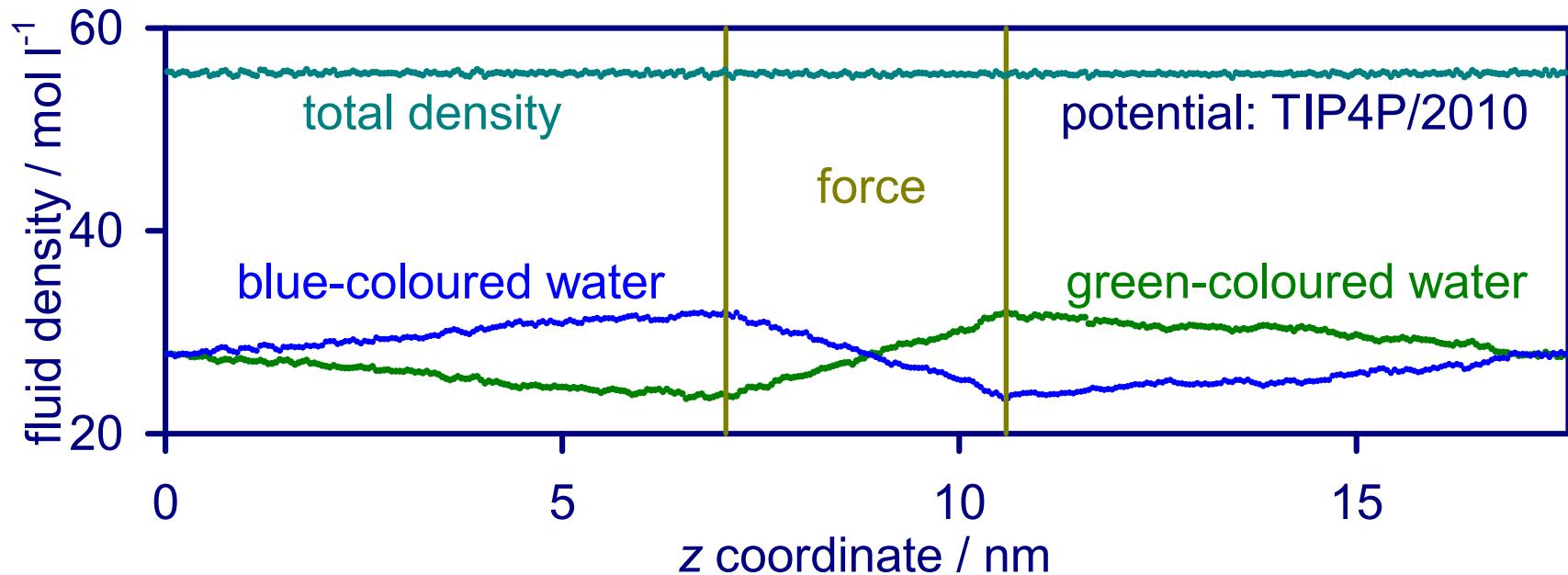
In principle, equilibrium MD computes  $D_s$ , whereas NEMD yields  $D_t$ .

**Note:** For a pure substance,  $\mu_i$  and  $p$  cannot be varied independently.



# Avendaño's dæmon

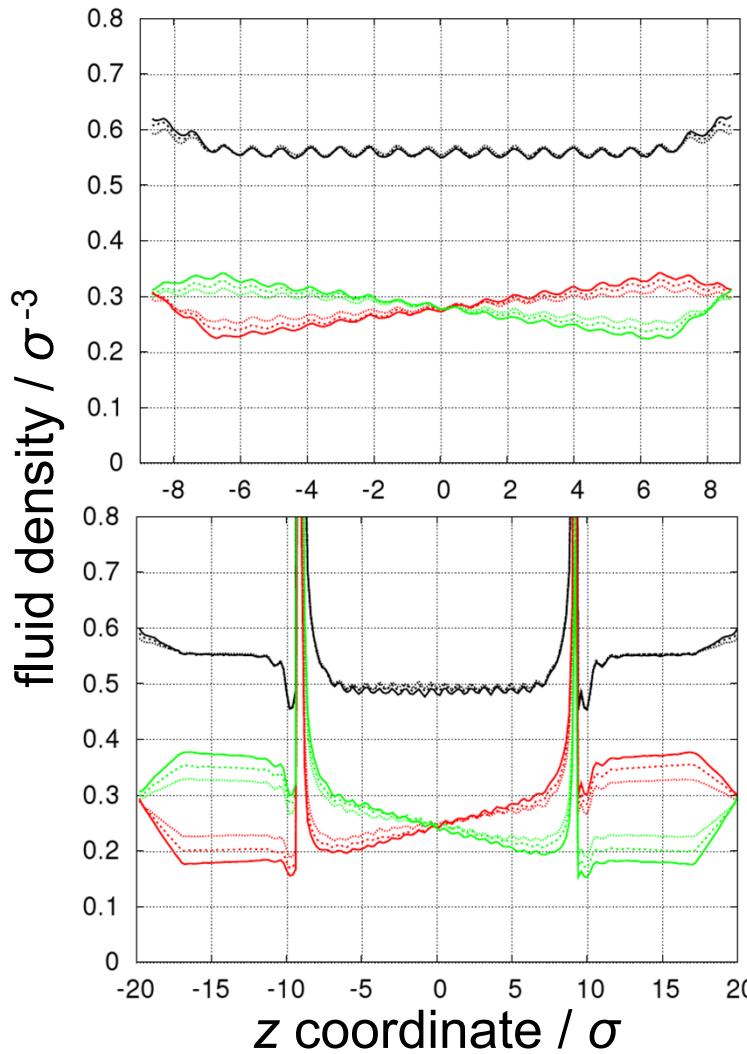
The mobility  $D_0$  can be determined by NEMD if a gradient in  $\mu$  ( $\rightarrow$  diffusive transport) is present without a pressure gradient ( $\rightarrow$  viscous transport).



Actually identical fluid molecules are assigned different labels (blue or green “colour”) and accelerated in opposite directions by Avendaño’s dæmon.



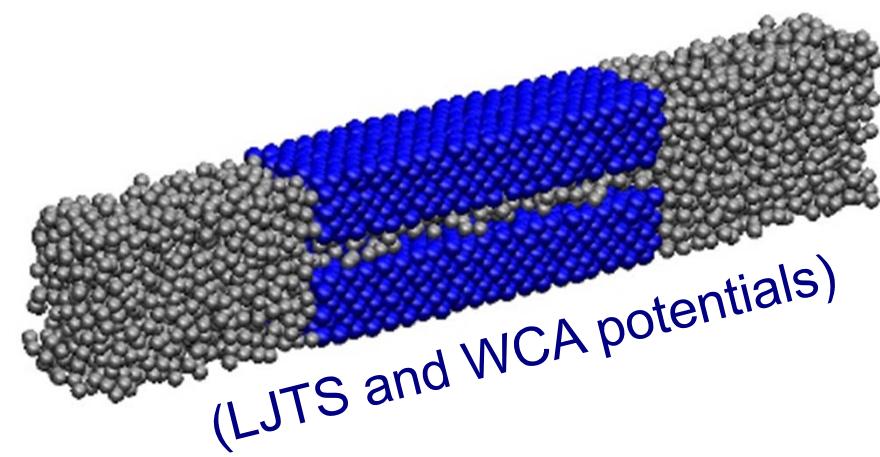
# Membrane topology



By comparison between the flow in

- an infinite (i.e. periodic) channel
- an “open” system including a bulk section ...

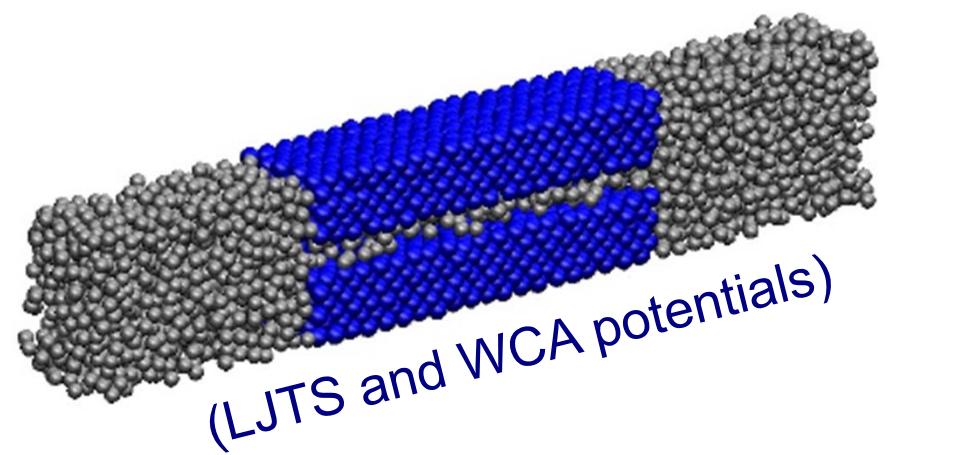
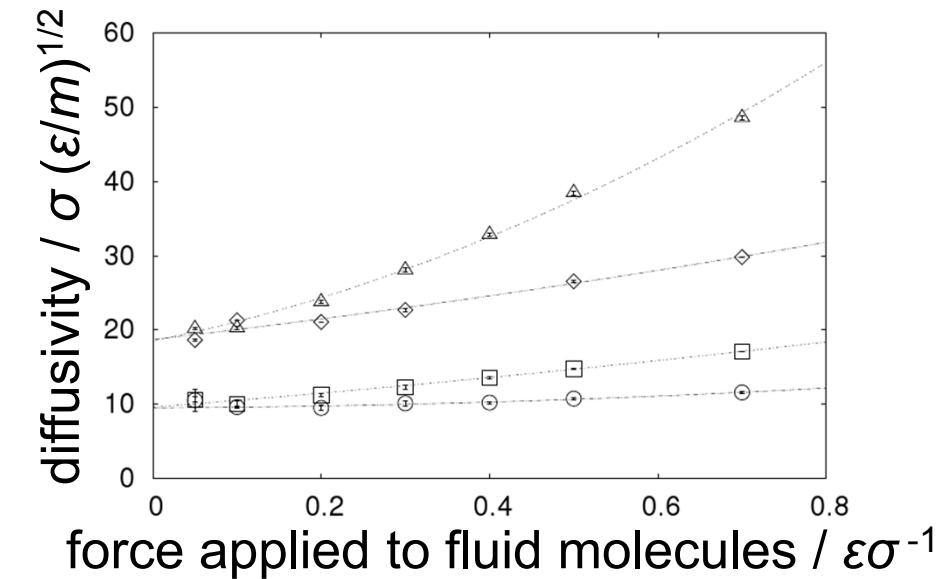
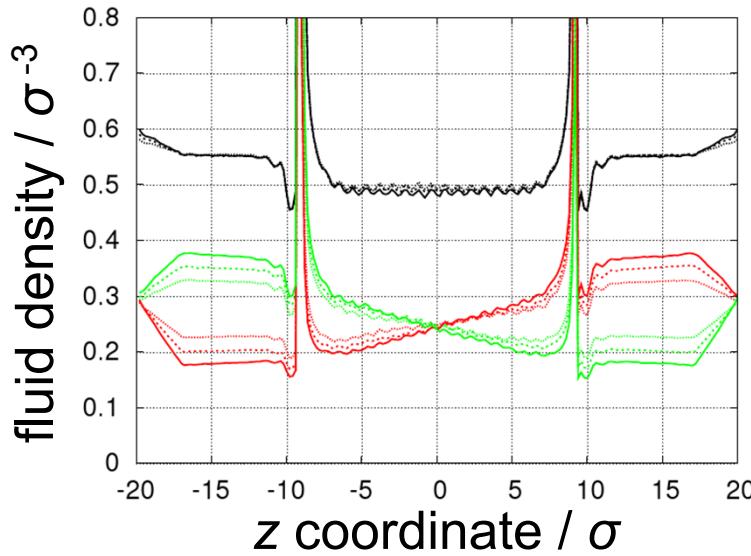
... the influence of the channel entrance and exit regions on the overall effective diffusivity can be isolated.





# Effective diffusivity

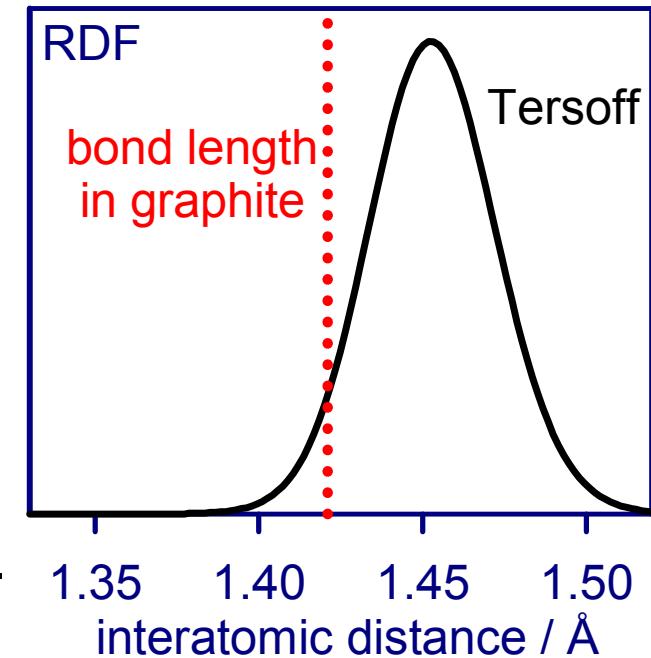
Diffusivities are obtained as linear response coefficients, i.e. in the limit where the accelerating force, which perturbs the equilibrium state of the system, approaches zero





# Nanofiltration membranes: Molecular model

Material	Chemical Structure	$T_g$ (°C)	Density (g/cm³)
Matrimid® 5218		323	1.22
P84		315	1.31



Tin et al. (2004), Ind. Eng. Chem. Res. **43**: 6467.

Length scale of the Tersoff potential for carbon:

Cutoff

$$R = 1.8 \text{ \AA}$$

$$S = 2.1 \text{ \AA}$$

Attraction

$$\mu = 2.2119 \text{ \AA}^{-1}$$

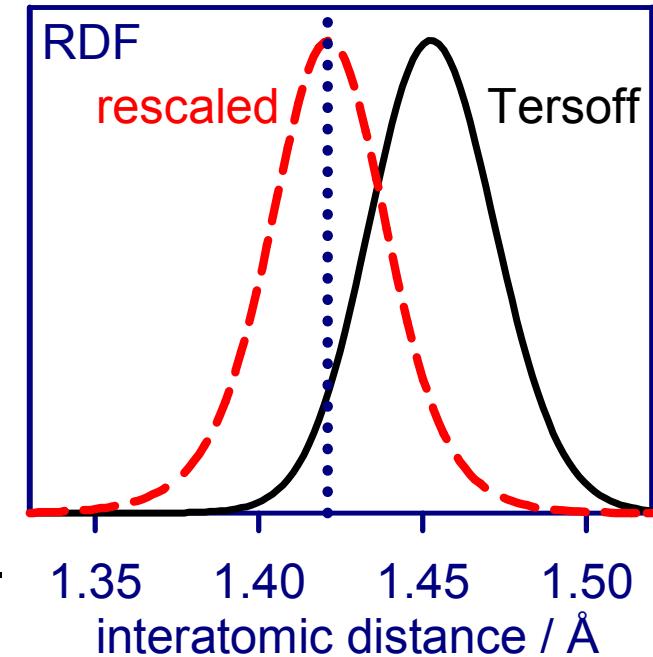
Repulsion

$$\lambda = 3.4879 \text{ \AA}^{-1}$$



# Nanofiltration membranes: Molecular model

Material	Chemical Structure	$T_g$ (°C)	Density (g/cm³)
Matrimid® 5218		323	1.22
P84		315	1.31



Tin et al. (2004), Ind. Eng. Chem. Res. **43**: 6467.

Optimized potential parameters for graphite:

Cutoff

$$R = 2.0 \text{ \AA} \quad (1.8 \text{ \AA})$$

$$S = 2.35 \text{ \AA} \quad (2.1 \text{ \AA})$$

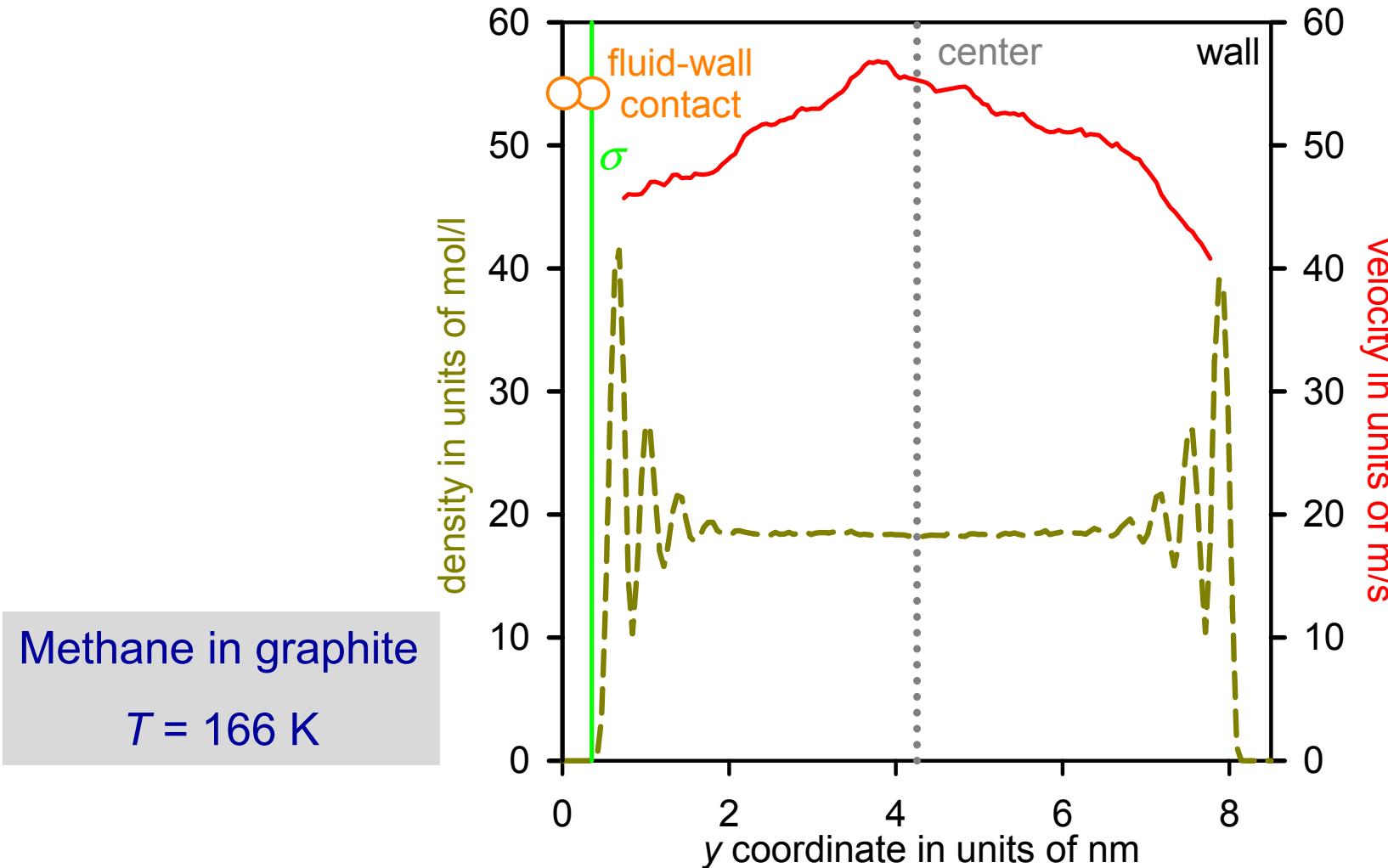
Attraction

$$\mu = 2.275 \text{ \AA}^{-1} \quad (2.2119 \text{ \AA}^{-1})$$

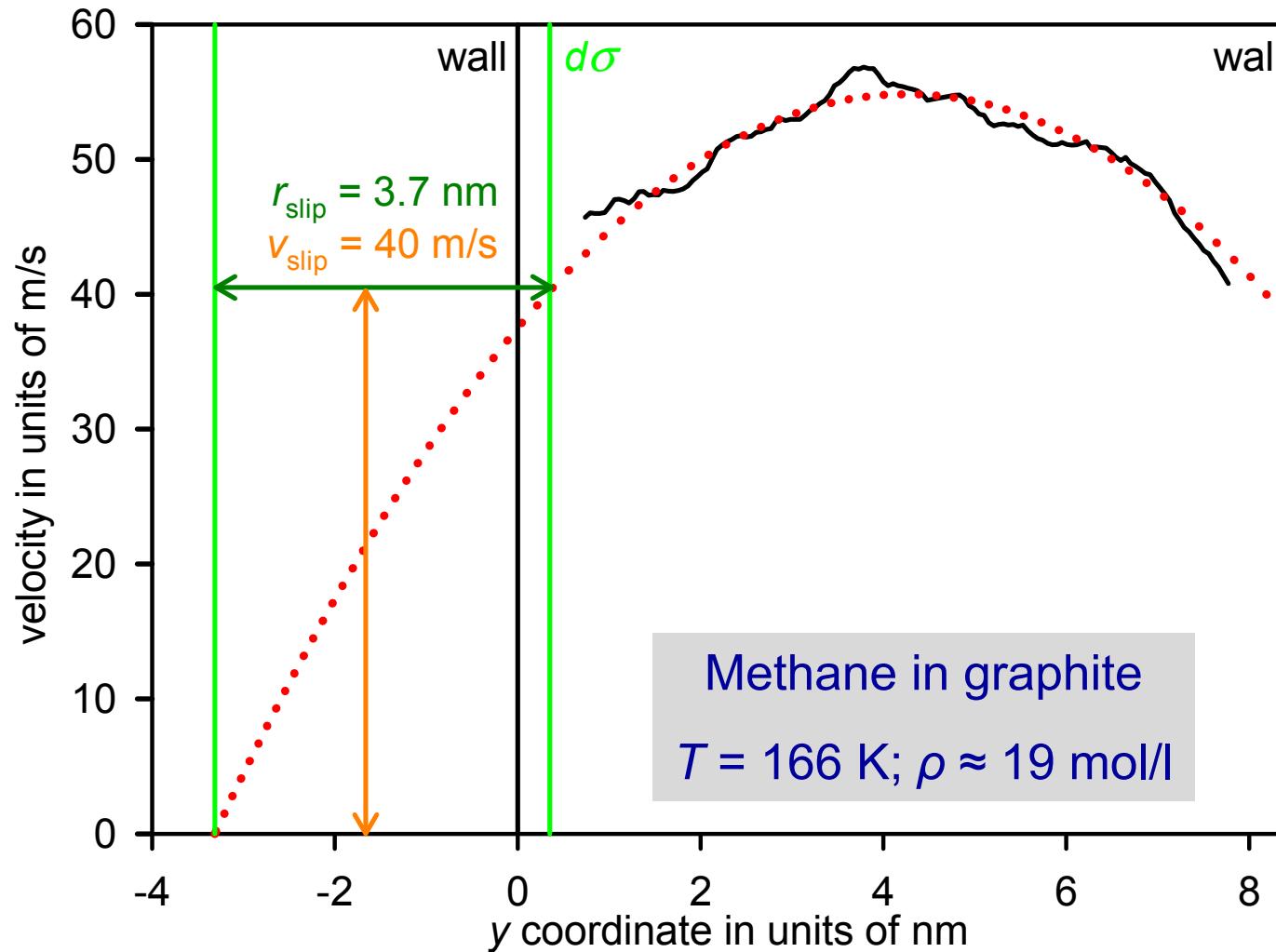
Repulsion

$$\lambda = 3.587 \text{ \AA}^{-1} \quad (3.4879 \text{ \AA}^{-1})$$

# Velocity profile and boundary slip



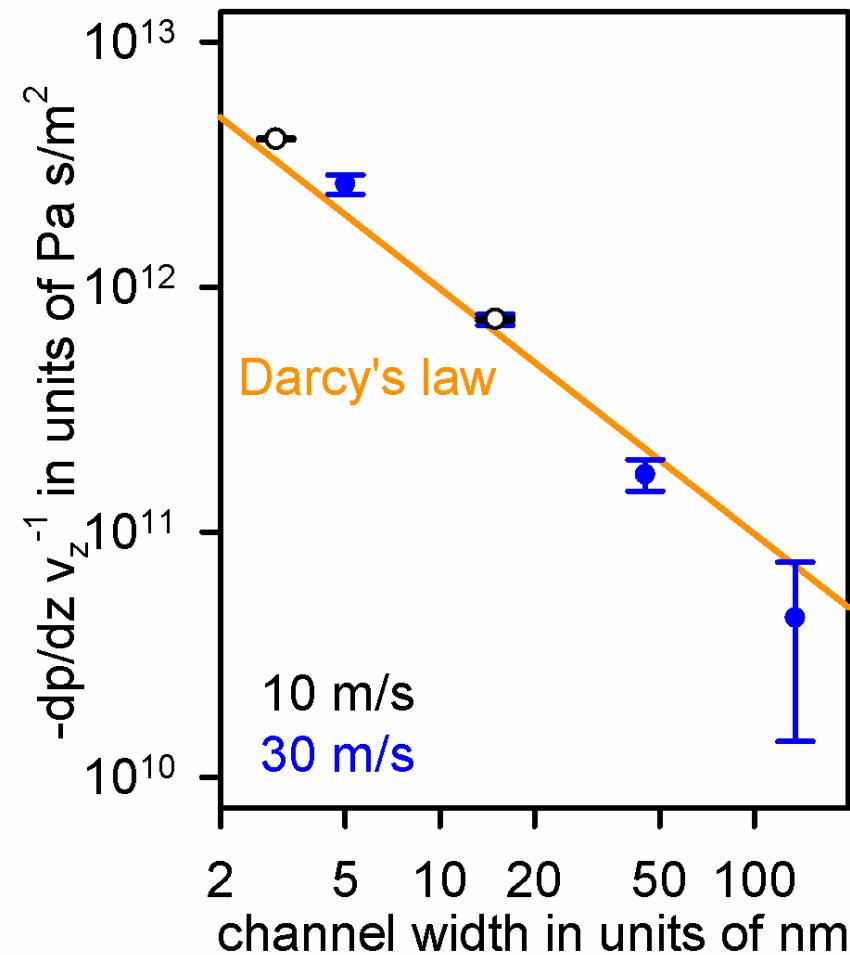
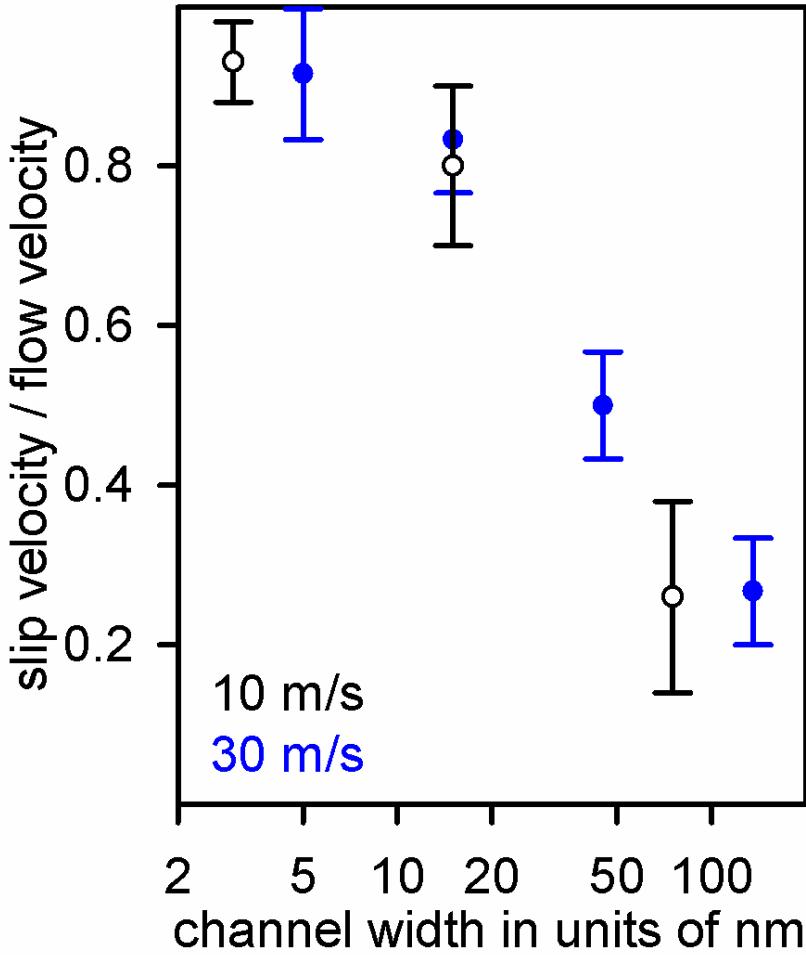
# Velocity profile and boundary slip





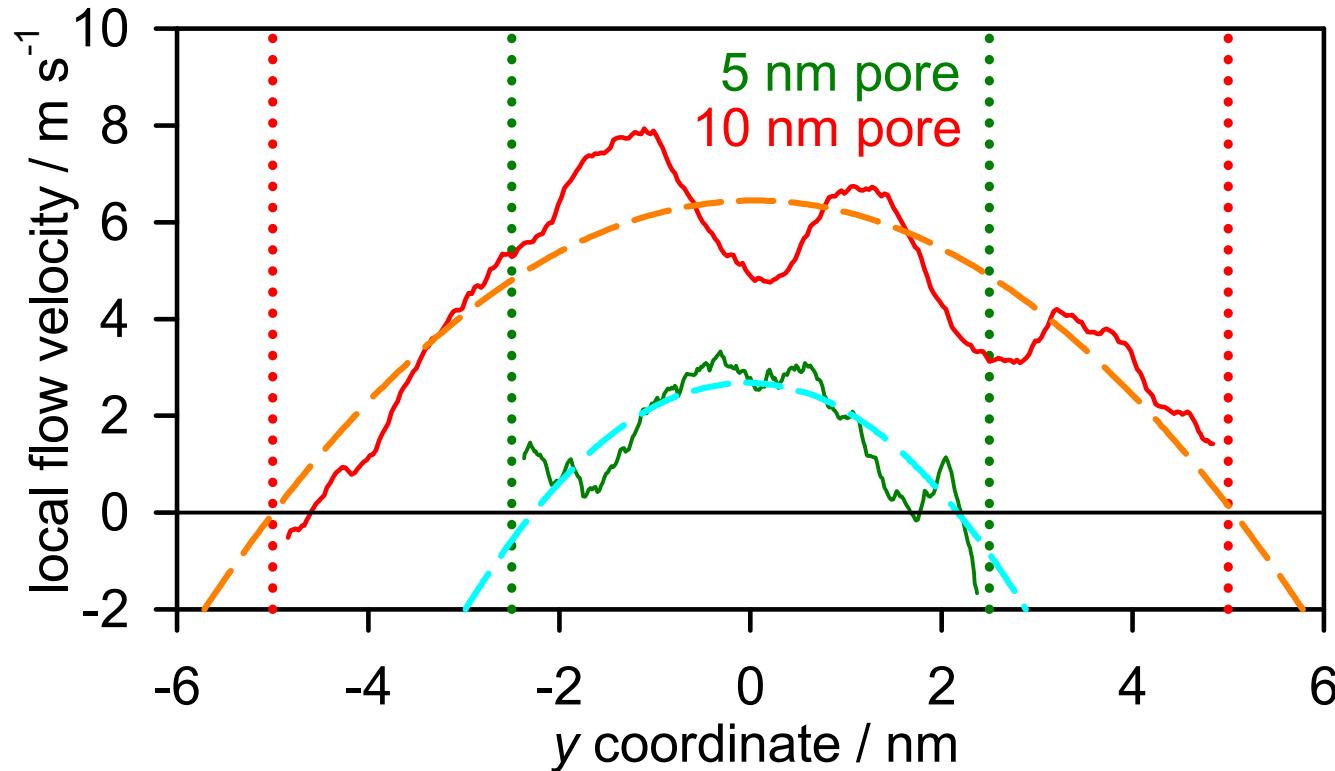
# From nanofluidics to microfluidics

Methane in graphite:  $T = 166$  K; values of  $\eta$  and  $\xi$  from Wang et al.





# Water in polar nanofiltration membranes



boundary slip  
practically  
negligible, in  
contrast to the  
unpolar system

$r_{\text{slip}}$  approaches  
the nm length  
scale at high  
flow velocities  
(200 m/s)

Quantitative water model:  
TIP4P/2010 (Huang *et al.*)

Qualitative P84 polyimide pore model:  
Graphite with superimposed point charges



# Massively parallel MD simulation

MD code ls1 mardyn (“large systems 1: molecular dynamics”)

## spatial domain decomposition

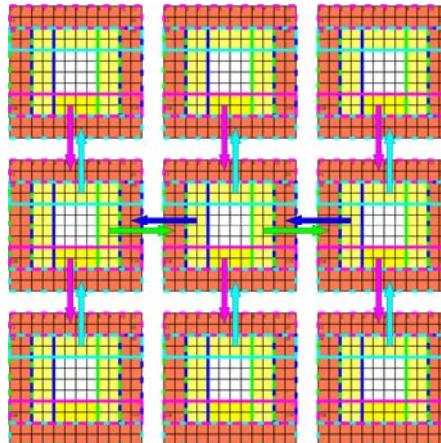
exploits the concurrency due to  
the limited range of the interactions

central, marginal, and halo cells

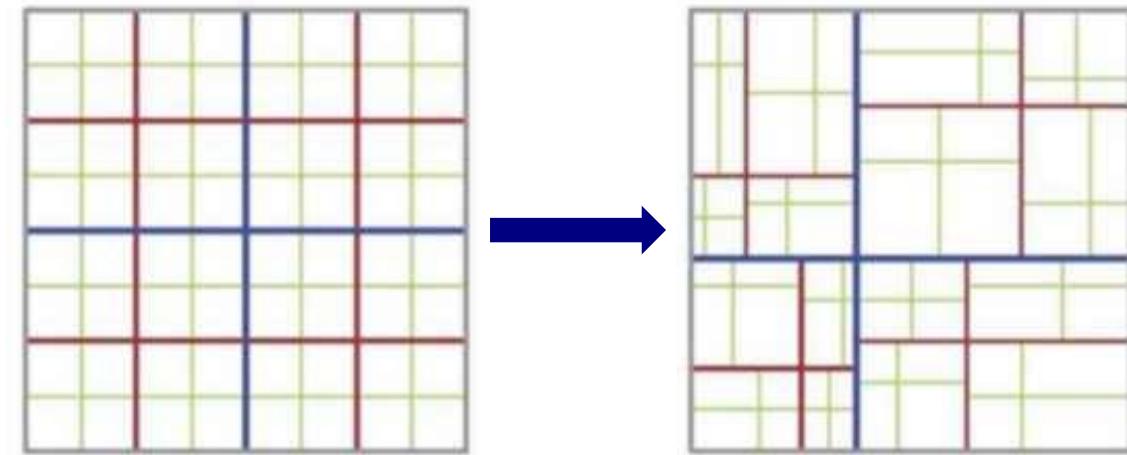
## (dynamic) load balancing

the simulated system and the HPC  
hardware may be heterogeneous

uses octrees or space-filling curves



linked-cell algorithm

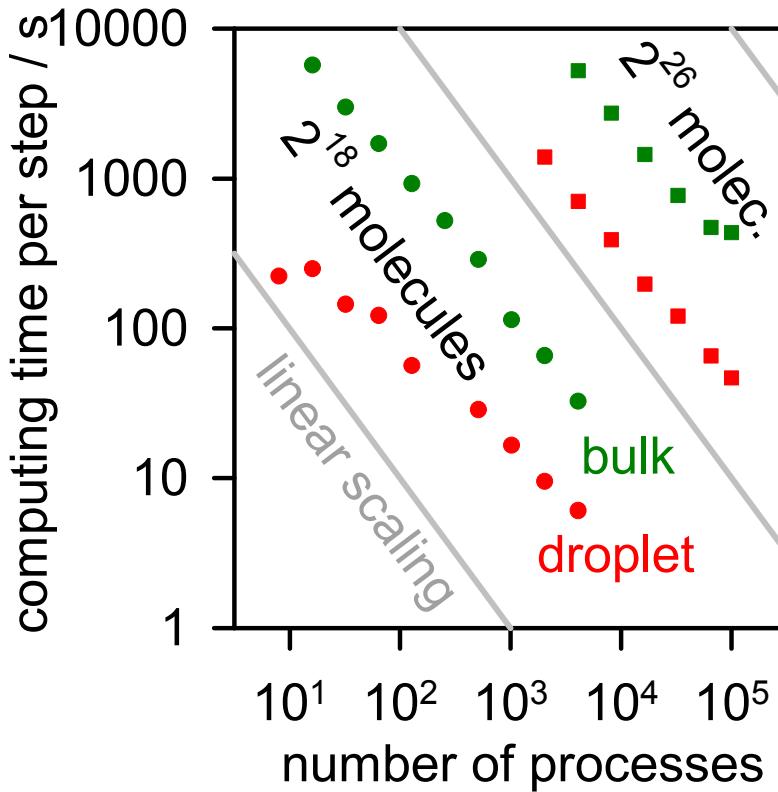


load balancing based on  $k$ -dimensional trees

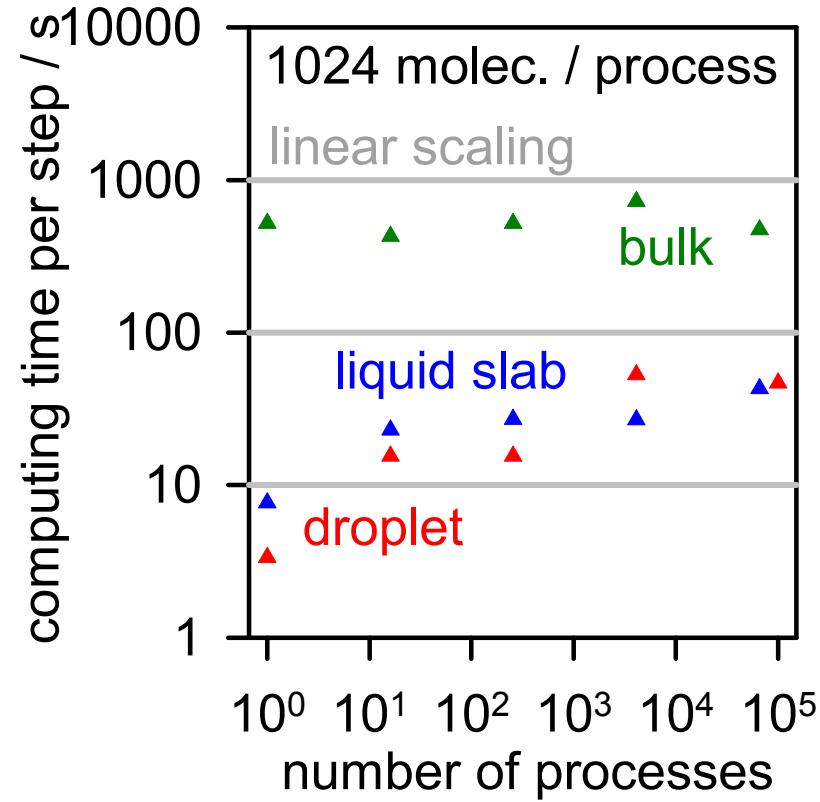
# Massively parallel MD simulation

MD code ls1 mardyn: scaling on the *Hermit* (top 12) at HLRS

**strong scaling (Amdahl)**



**weak scaling (Gustafson)**





# Conclusion

- ⇒ Poiseuille flow can be investigated by non-equilibrium MD simulation, i.e., by compensating or overcompensating the pressure drop.
- ⇒ Avendaño's dæmon makes purely diffusive transport (i.e., the mobility coefficient) accessible to NEMD for confined systems.
- ⇒ For methane in graphite, Darcy's law was found to hold down to the molecular length scale significant boundary slip was present for diameters below 100 nm.
- ⇒ In case of water in a polar membrane material, no significant boundary slip was detected.
- ⇒ Massively parallel MD (e.g. with ls1 mardyn) promises to make a molecular analysis of microfluidics feasible within the present decade.