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Role of the frame of reference in the apparent kinetic contribution to the surface tension

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1. <u>The established view</u>

- 2. Apparent kinetic contribution
- 3. Interfacial properties: Open questions

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Surface tension: Macromechanics



- Mechanical equilibrium \rightarrow fluid at rest, off-diagonal entries in pressure tensor are zero.
- The tangential pressure $P_{t} = P^{xx} = P^{zz}$ deviates from the normal pressure $P_{n} = P^{yy}$.
- Compared to the system boundaries perpendicular to *y* direction (normal direction), those in tangential direction experience an additional tangential force inward:
 - In x direction, $f_x \sim L_z$; similarly, in z direction, $f_z \sim L_x$.
 - By symmetry, the proportionality factors are equal, $f_x = -\gamma L_z$ and $f_z = -\gamma L_x$.
- For x direction, $P^{xx} P_n = f_x / L_y L_z = -\gamma L_z / L_y L_z = -\gamma / L_y$. Therefore, $\gamma = L_y (P_n P^{xx})$.
- The same argument applies in z direction. So we can write $\gamma = L_y (P_n P_t)$.

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From macro- to micromechanics

Using localized **pressure profiles** based on **continuum mechanics**:

 $y = \int_{y} dy \left(P_{n}(y) - P_{t}(y) \right)$

vapour-liquid surface tension from integral over deviation between P_n and P_t

Compare the expression using the pressure P_n and P_t for the whole system:

- For x direction, $P^{xx} P_n = f_x / L_y L_z = -\gamma L_z / L_y L_z = -\gamma / L_y$. Therefore, $\gamma = L_y (P_n P^{xx})$.
- The same argument applies in z direction. So we can write $\gamma = L_y (P_n P_t)$.

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Grounding in statistical mechanics

Using localized **pressure profiles** based on **continuum mechanics**:

$$y = \int_{y} dy \left(P_{n}(y) - P_{t}(y) \right)$$

In a **particle-based system**, localized (continuum) pressures are **not uniquely defined**. But their volume integral is uniquely defined, and therefore γ is as well – in the case of a planar interface.



392: 2359, **2013**.

Pressure obtained as $P_n = \rho T + V^{-1} < \Pi_n > \text{ and } P_t = \rho T + V^{-1} < \Pi_t >$.

- Therein, Π is the virial, $\Pi_n = \Pi^{yy} = \sum_{\{i,j\}} r_{ij,y} f_{ij,y}$ and $\Pi_t = \Pi^{xx} = \sum_{\{i,j\}} r_{ij,x} f_{ij,x}$.
- Since $V = aL_y$, where a is the surface area, $\gamma = a^{-1} (\langle \Pi_n \rangle \langle \Pi_t \rangle)$. Note that ρT cancels out!
- While the virial for the whole system is uniquely defined, the local profiles are not.

Grounding in statistical mechanics

One way of obtaining a statistical mechanical expression for the surface tension is through a hypothetical **infinitesimal test transformation**:



canonical ensemble

- **N**, *V* and *T* constant
- V' and V'' (liq, vap) constant
- Surface area *a* changes

$$\gamma = \left(\frac{\partial F}{\partial a}\right)_{\mathbf{N}, V', V'', T}$$

For a system of point masses, this leads to the **same result as the argument** from mechanics, $\gamma = a^{-1}(\langle \Pi_n \rangle - \langle \Pi_t \rangle)$; there is no contribution from ρT .

Pressure obtained as $P_n = \rho T + V^{-1} < \Pi_n > \text{ and } P_t = \rho T + V^{-1} < \Pi_t >$.

- Therein, Π is the virial, $\Pi_n = \Pi^{yy} = \sum_{\{i,j\}} r_{ij,y} f_{ij,y}$ and $\Pi_t = \Pi^{xx} = \sum_{\{i,j\}} r_{ij,x} f_{ij,x}$.
- Since $V = aL_{y}$, where a is the surface area, $\gamma = a^{-1} (\langle \Pi_n \rangle \langle \Pi_t \rangle)$. Note that ρT cancels out!
- While the virial for the whole system is uniquely defined, the local profiles are not.

Rigid multisite models: Surface tension



Systematic exploration of the physically relevant part of the model parameter space
Correlation of the 2LJCQ and 2LJCD surface tension by critical-scaling expressions

¹S. Werth et al., Chem. Eng. Sci. **121**: 110-117, doi:10.1016/j.ces.2014.08.035, **2015**.

²S. Werth, M. Horsch, H. Hasse, J. Chem. Phys. **144**: 054702, doi:10.1063/1.4940966, **2016**.

Accuracy of surface tension from molecular models



¹S. Werth et al., Chem. Eng. Sci. **121**: 110-117, doi:10.1016/j.ces.2014.08.035, **2015**.

²S. Werth, M. Horsch, H. Hasse, J. Chem. Phys. **144**: 054702, doi:10.1063/1.4940966, **2016**.

Deviation between surface

Accuracy of surface tension from molecular models

Models' overestimation of γ by different model classes on average:

20 %

Dipolar: 2CLJD (contd.)

R30B2 (CH_2Br_2) Methylene iodide (CH₂I₂) R12B2 (CBr_2F_2) R12B1 (CBrClF₂) R10B1 (CBrCl₂) R161 (CH,F-CH,) R150a (CHCl₂-CH₂) R140 (CHCl₂-CH₂Cl) R140a (CCl₃-CH₃) R130a (CH₂CI-CCl₂) R160B1 (CH₂Br-CH₃) R150B2 (CHBr₂-CH₂) R131b (CH₂F-CCl₂) R123B1 (CHClBr-CF_a) R112a (CCl₃-CF₂Cl) 12 % R1141 (CHF=CH₂) R1132a (CF₂=CH₂) R1140 (CHCI=CH₂) R1122 (CHCI=CF_a) R1113 (CFCI=CF.) R1113B1 (CFBr=CF₂)

Dipolar: 2CLJD

R11 (CFCl₂)

R12 (CF₂Cl₂)

R13 (CF₃Cl)

R13B1 (CBrF₃)

R22 (CHF₂CI)

R123 (CHCl₂-CF₂)

R124 (CHFCI-CF₃)

R125 (CHF₂-CF₃)

R134a (CH₂F-CF₂)

R141b (CH₃-CFCl₂)

R142b (CH₂-CF₂Cl)

R143a (CH₃-CF₃)

R40 (CH₂Cl)

R20 (CHCl₃)

R20B3 (CHBr₂)

R21 (CHFCl₂)

R32 (CH₂F₂)

R30 (CH₂Cl₂)

R40B1 (CH₂Br)

R152a (CH₂-CHF₂)

Methyl iodide (CH₃I)

R30B1 (CH₂BrCl)

R23 (CHF_a)

R41 (CH₃F)

Carbon monoxide (CO)

Quadrupolar: 2CLJQ

Fluorine (F_2) Chlorine (Cl_2) Bromine (Br_2) Iodine (l_2) Nitrogen (N_2) Oxygen (O_2) Carbon dioxide (CO_2)

Quadrupolar: 2CLJQ (contd.)

Carbon disulfide (CS₂) Ethane (C_2H_6) Ethylene (C_2H_4) Acetylene (C₂H₂) R116 (C_2F_6) R1114 (C₂F₄) R1110 (C_2CI_4) Propadiene (CH₂=C=CH₂) Propyne (CH_3 - $C\equiv CH$) Propylene (CH₂-CH=CH₂) R846 (SF_e) R14 (CF₄) R10 (CCl₄) R113 (CFCl,-CF,Cl) R114 (CF₂CI-CF₂CI) 22 % R115 (CF₂-CF₂Cl) R134 (CHF₂-CHF₂) R150B2 (CH₂Br-CH₂Br) R114B2 (CBrF₂-CBrF₂) R1120 (CHCI=CCI_a)

Multi-site models

Isobutane (C_4H_{10}) Cyclohexane (C_6H_{12}) Methanol (CH₃OH) Ethanol (C₂H₅OH) Formaldehyde (CH₂=O) Dimethylether (CH₃-O-CH₃) Acetone (C₃H₆O)

Multi-site models (contd.)

Ammonia (NH₂) Methylamine (NH₂-CH₂) Dimethylamine (CH₂-NH-CH₂) R227ea (CF₃-CHF-CF₃) Sulfur dioxide (SO₂) Ethylene oxide (C_2H_4O) Dimethylsulfide (CH₃-S-CH₃) Hvdrogen cyanide (HCN) Acetonitrile (NC₂H₃) Thiophene (SC_4H_4) Nitromethane $(NO_{2}CH_{2})$ Phosgene (COCl₂) Benzene ($C_{e}H_{e}$) Toluene (C₂H₂) Chlorobenzene (C₆H₅Cl) Dichlorobenzene ($C_6H_4CI_2$) Cyclohexanol (C₆H₁₁OH) Cyclohexanone ($C_{e}H_{10}O$) Cvanogen (C₂N₂) Cyanogen chloride (CCIN) Formic acide (CH_2O_2) Ethylene glycol ($C_2H_6O_2$) Water (H₂O) Hydrazine (N_2H_4) Methylhydrazine (CH_eN₂) Dimethylhydrazine ($C_2H_8N_2$) Perfluorobutane (C₄F₁₀) Ethyl acetate $(C_4H_9O_2)$ Hexamethyldisiloxane (C₆H₁₂OSi₂) Octamethylcyclotetrasiloxane (C8H24O4Si4)

If the employed method was wrong, all this work (and more) must be redone!
 If the models' actual surface tension is greater, the models are even worse!



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The apparent kinetic contribution¹

At interfaces, there can be a **preferred orientation**:



Sega *et al*.^{1, 2} distinguish:

surface tension γ^{Ξ} computed using the **virial tensor only** (which, they claim, is inaccurate) surface tension γ^{p} computed using the **micromechanical pressure** tensor Assume, e.g., most molecules are arranged normal to the interface, where the models are rigid. The **rotation** then contributes much more to **motion in tangential direction** than normal direction.

$$p_{ab} = \frac{2}{V} \left(\frac{1}{2} \sum_{i} m_i v_i^a v_i^b - \Xi_{ab} \right)$$

micromechanical pressure following Sega *et al*.²

If, as they claim, this is the right definition:

- all our simulations were wrong
- models might become worse

¹M. Sega, B. Fábián, P. Jedlovszky, *J. Phys. Chem. Lett.* **8**: 2608-2612, doi:10.1021/acs.jpclett.7b01024, **2017**. ²M. Sega, G. Horvai, P. Jedlovszky, *J. Mol. Liq.* **262**: 58-62, doi:10.1016/j.molliq.2018.04.004, **2018**.

The argument from intuition

Case 1: Tangential orientation preferred.



Rotation contributes most to *normal motion* of sites.

Kinetic energy and micromechanical pressure in tangential direction are lowest.

This works with surface tension, so that γ is increased.

Case 2: Normal orientation preferred.



Rotation contributes most to *tangential motion* of sites.



This works against surface tension, so that γ is reduced.

The argument¹ according to Lbadaoui-Darvas et al.²

The argument is presented by M. Lbadaoui-Darvas et al.² as follows:

"Although the virial route has the advantage that it only requires molecular coordinates [...] while for the pressure route velocities are also needed, the underlying assumption is not always true [...]. Namely, the **equipartition theorem states that** the average of **the total energy** of the particles (i.e., the sum of their potential and kinetic energies) rather than the kinetic energy itself **is distributed evenly along** all **spatial directions**."

$$p_{\alpha\beta} = \frac{1}{V} \left(\sum_{i} m_{i} v_{i}^{\alpha} v_{i}^{\beta} + \left\langle \sum_{ij} f_{ij}^{\alpha} \int_{C_{ij}} \delta(\mathbf{r} - \mathbf{s}) \, \mathrm{d}s^{\beta} \right\rangle \right)$$
(3)

Sega *et al.*¹ distinguish: surface tension γ^{Ξ} which, using the **statistical-mechanical** pressure, is $L_y(P_n - P_t)$ surface tension γ^p which, using the **micromechanical** pressure, is $L_y(P_n - P_t)$

pressure tensor following Lbadaoui-Darvas *et al*.²

(just another notation for the **micromechanical pressure tensor** observable)

¹M. Sega, B. Fábián, P. Jedlovszky, J. Phys. Chem. Lett. 8: 2608-2612, doi:10.1021/acs.jpclett.7b01024, 2017.
 ²M. Lbadaoui-Darvas et al., J. Phys. Chem. B 126(4): 751-765, doi:10.1021/acs.jpcb.1c08553, 2022.
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Magnitude of the apparent kinetic contribution¹



deviation between P_n and P_t

Figure 3. Upper panel: surface tension γ^p and virial contribution γ^{Ξ} as a function of temperature. Lower panel: ratio $\gamma^{id}/\gamma^p = 1 - \gamma^{\Xi}/\gamma^p$. Error bars are always smaller than the symbols in the upper panel and are for all temperatures of the order of 0.1–0.2 mN/m.

¹M. Sega, B. Fábián, P. Jedlovszky, J. Phys. Chem. Lett. 8: 2608-2612, doi:10.1021/acs.jpclett.7b01024, 2017.

Magnitude of the apparent kinetic contribution¹





order parameter^{2, 3} <3 cos² θ – 1> / 2 (evaporation, not equilibrium)

For SPC/E water, Sega *et al.*¹ obtain a surface tension γ^p that is increased by 16% over γ^{Ξ} .

Orientation profiles for Stockmayer and 2CLJD done by Mecke *et al.*² long ago.

Results from Homes *et al.*³ on 2CLJ and 2CLJQ.

¹M. Sega, B. Fábián, P. Jedlovszky, J. Phys. Chem. Lett. 8: 2608-2612, doi:10.1021/acs.jpclett.7b01024, 2017.
 ²M. Mecke, J. Fischer, J. Winkelmann, J. Chem. Phys. 114(13): 5842-5852, doi:10.1063/1.1349177, 2001.
 ³S. Homes, M. Heinen, J. Vrabec, Phys. Fluids 35: 052111, doi:10.1063/5.0147306, 2023.

Pressure-tensor observables

From work preceding Sega *et al.*,¹ two different expressions were often used for the *microscopic pressure-tensor observable* in the literature:

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statistical-mechanicalmicromechanical
$$P = 1\rho T + V^{-1} < \Pi >$$
 $P = V^{-1} < K + \Pi >$ where $\Pi = \sum_{i,j} \mathbf{r}_{ij} \otimes \mathbf{f}_{ij}$ where $\Pi = \sum_{i,j} \mathbf{r}_{ij} \otimes \mathbf{f}_{ij}$ and $K = \sum_{i} m_{i} \mathbf{v}_{i} \otimes \mathbf{v}_{i}$

The two definitions were thought of as equivalent. The merit of the Sega *et al.*¹ paper consists in showing that, for rigid multi-site models, there is a deviation.

¹M. Sega, B. Fábián, P. Jedlovszky, *J. Phys. Chem. Lett.* **8**: 2608–2612, doi:10.1021/acs.jpclett.7b01024, **2017**. CECAM IPOQ 2025 30th January 2025 16

Pressure-tensor observables

From work preceding Sega *et al.*,¹ two different expressions were often used for the *microscopic pressure-tensor observable* in the literature:



The two definitions were thought of as equivalent. The merit of the Sega *et al.*¹ paper consists in showing that, for rigid multi-site models, there is a deviation if the sum is evaluated over coordinates of individual interaction sites / atoms.

However, note that in our statistical-mechanical expressions, the momentum coordinates refer to *molecular centres of mass* as well as *angular momenta*.

¹M. Sega, B. Fábián, P. Jedlovszky, J. Phys. Chem. Lett. 8: 2608-2612, doi:10.1021/acs.jpclett.7b01024, 2017.CECAM IPOQ 202530th January 202517



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What does this all mean for us?

The micromechanical pressure tensor is a standard literature expression, going back to Irving and Kirkwood.¹ It is the one implemented in GROMACS.²

But our pressure and surface tension expressions are strictly the ones from statistical mechanics. We don't need to add an apparent kinetic contribution. But if we do it consistently over rigid units, it should be zero. The micro-mechanical pressure could be helpful for connecting to continuum mechanics.

There are two possibilities regarding codes using SHAKE, RATTLE, or similar:

- The apparent kinetic contribution is really needed if rigid molecular models are analysed over effective forces acting on individual atoms. *This is possible, but a clear proof of it cannot be found in the literature.*
- 2) But it is entirely possible that even then, the real γ is γ^{Ξ} , and not γ^{p} . While $\gamma^{\Xi} \neq \gamma^{p}$ was shown, the literature contains no clear proof of $\gamma = \gamma^{p}$.

¹J. H. Irving, J. G. Kirkwood, *J. Chem. Phys.* **18(6)**: 817–829, **1950**; therein, see Eqs. (5.12) to (5.14). ²GROMACS manual, https://manual.gromacs.org/, section on "the global MD algorithm," **2023**.

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